

Using cgs units we get:

$$v_t = (1 \text{ ft sec}^{-2})(12 \text{ in ft}^{-1})(2.54 \text{ cm in}^{-1}) = 30.5 \text{ cm sec}^{-1}$$

$$\rho = (0.045 \text{ lb}_m \text{ ft}^{-3})(454 \text{ g lb}_m^{-1})(12 \times 2.54)^{-3} (\text{ft}^3 \text{ cm}^{-3})$$

$$= 7.2 \times 10^{-4} \text{ g cm}^{-3}$$

Therefore:

$$D_{\max} = \sqrt{\frac{18 \mu v_t}{(\rho_s - \rho) g}} = \sqrt{\frac{(18)(0.00026)(30.5)}{(1.2 - 7.2 \times 10^{-4})(981)}}$$

$$= \sqrt{1.21 \times 10^{-4}} = 1.1 \times 10^{-2} \text{ cm} = 110 \text{ microns}$$

(b) Eq. 2.6-16 is of course valid only for $Re < 0.1$ although it may be applied approximately up to about $Re = 1$. For the system at hand

$$Re = \frac{D v_t \rho}{\mu} = \frac{(1.1 \times 10^{-2})(30.5)(7.2 \times 10^{-4})}{(0.00026)} = 0.95$$

In Chapter 6, methods are given for handling flow around spheres for $Re > 1$.

2.D Flow of Falling Film--- Alternate Derivations

(a) Set up a momentum balance as before and obtain the differential equation:

$$\frac{d\tau_{xz}}{dz} = \rho g \cos \beta$$

which may be integrated to give:

$$\tau_{xz} = \rho g z \cos \beta + C_1$$

Inasmuch as no momentum is transferred at $z = \delta$, then at that plane we have $\tau_{xz} = 0$. This boundary condition enables us to determine C_1 to be:

$$C_1 = -\rho g \delta \cos \beta$$

and hence the momentum flux distribution is:

$$\tau_{xz} = -\rho g \delta \cos \beta [1 - (z/\delta)]$$

Note that the momentum flux is in the negative z -direction.

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R Bogdan



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